Spin state relaxation of iron complexes: the case for OPBE and S12g

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Abstract: The structures of nine iron complexes that show a diversity of experimentally 13 observed spin ground states are optimized and analysed with Density Functional Theory 14 (DFT). An extensive validation study of the new S12g functional is performed, with the 15 discussion concerning the influence of the environment, geometry and its overall 16 performance based on the comparison with the well proven OPBE functional. The OPBE and 17 18 S12g functionals give the correct spin ground state for all investigated iron complexes. Since S12g performs remarkably well it can be considered a reliable tool for studying spin state 19 energetics in complicated transition metal systems. 20

21 Keywords: Density Functional Theory; Fe(II) and Fe(III) coordination compounds; validation study; spin states. 22

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RUNNING TITLE: SPIN STATE ENERGETICS OF IRON COMPLEXES 24

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INTRODUCTION

Spin is an intrinsic and inherent property of atoms and molecules.¹ Most of the transition 27 metal ions with partially filled d-shells can exhibit different kinds of spin multiplicity in the 28 ground state, i.e. can lead to different spin states. Depending on the oxidation number, iron 29 complexes have usually either 5 of 6 d-electrons that can be distributed in octahedral 30 environment in at least two different ways: with a maximum number of unpaired electrons, 31 leading to the high spin (HS) state, or with maximally paired electrons – giving the low spin 32 (LS) state. Other possibilities of the distribution of electrons represent the intermediate (IS) 33 spin state. Since HS, IS and LS complexes usually display quite different structural, spectral 34

and magnetic properties, and often also reactivity, it is of the utmost importance to have both experimental and theoretical methods to correctly determine the spin ground state of the system. However, both experiment and theory have difficulties and problems, and many studies have been devoted to this issue in the last decade.^{2, 3, 4, 5}

From a broad palette of quantum mechanical methods, Density Functional Theory (DFT)⁶, 39 ^{7, 8} emerged into the mainstream, mainly because it gives a good compromise between the 40 accuracy of the results and the computational efficiency.^{9, 10, 11} However, although DFT gives 41 in principle exact energy, the universal functional is still unknown, leading to density 42 43 functional aproximations (DFAs). These DFAs are parameterized for different properties and, noteworthy, spin-state energies were not included in the development for most of nowadays 44 available DFAs.¹² It has been shown that the accuracy of the results strongly depends of the 45 choice of the DFAs, but also on the basis set that is used.^{1, 3, 13, 14} Early pure functionals like 46 LDA^{15, 16, 17}, BP86,^{18, 19} BLYP,^{19, 20} or PW91^{21, 22} have a tendency to favor LS states,¹⁴ while 47 hybrid functionals like B3LYP^{23, 24}, PBE0²⁵ and M06^{26, 27} systematically favor HS.¹⁴ For the 48 reliable prediction of the correct spin ground state from a number of close lying states, 49 OPBE¹⁴ emerged to be one of the best functionals for the job.²⁸ Recently Swart constructed a 50 new density functional that combines the best of OPBE (spin states, reaction bariers) with the 51 best of PBE (week interactions) into the S12g⁵ DFA. 52

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Previously we reported the relative spin state energies of seven iron complexes (**1-7**, Figure 1) on OLYP^{20, 29} optimized geometries (**1-3**) and on crystal structures (**4-7**) with a variety of DFAs and already showed the good performance of OPBE for vertical spin state splittings.¹⁴

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- 59

<Figure 1>

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Herein, we present a detailed DFT study on OPBE optimized geometries of iron complexes (1-7) with experimentally established spin ground states, ranging from singlet to sextet, and extended it with two iron porphyrinato complexes (8, 9, Figure 1) that have been reported to have different electronic ground state in spite of their similarity.^{30, 31, 32}

Furthermore, a comprehensive validation study of the S12g DFA,⁵ together with the examination of the influence of the chemical environment is performed on all investigated complexes.

METHODOLOGY

70 All DFT calculations were performed with the Amsterdam Density Functional (ADF) suite of program.^{33,34} MOs were expanded in an uncontracted set of Slater type orbitals (STOs) of 71 triple- ζ quality containing diffuse functions (TZP)³⁵ and one set of polarization functions. 72 Core electrons (1s for 2nd period, 1s2s2p for 3rd-4th period) were not treated explicitly during 73 the geometry optimizations (frozen core approximation), as it was shown to have a negligible 74 effect on the obtained geometries.³⁶ An auxiliary set of s, p, d, f, and g STOs was used to fit 75 the molecular density and to represent the Coulomb and exchange potentials accurately for 76 each SCF cycle. 77

Energies and gradients were calculated using the OPBE and S12g functionals, in the gas phase and with the COSMO^{37, 38, 39} dielectric continuum model for the solvent environment. Geometries were optimized with the QUILD program⁴⁰ using adapted delocalized coordinates⁴¹ until the maximum gradient component was less than 10⁻⁴ a.u. Subsequent single point calculations that utilize all electron basis set have been performed on all optimized geometries, with OPBE and S12g.

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RESULTS AND DISCUSSION

- The total set of molecules consists both of Fe(III) (1-3, 8-9) and Fe(II) (4-7) complexes, 87 and show a diversity of experimentally observed spin ground states. A thorough examination 88 89 with the OPBE and S12g functionals, in the gas phase and the COSMO solvent environment, is performed. We will start our discussion focusing on the influence of the structure 90 relaxation on the spin states of Fe(III)-complexes 1-3.⁴² Experimentally, Fe-(PyPepS)₂ (1, 91 PyPepSH₂=N-2-mercaptophenyl-2'-pyridine-carboxamide) has a LS doublet ground state,⁴³ 92 Fe(tsalen)Cl (2, tsalen = N,N'-ethylenebis-(thio-salicylideneiminato)) an intermediate spin 93 (IS), quartet ground state⁴⁴ and $Fe(N(CH_2-o-C_6H_4S)_3)(1-Me-imidazole)$ (3) a HS sextet 94 ground state.⁴⁵ Then we will discuss the Fe(II)-complexes ((Fe(NH)S₄)L, (NH)S₄=bis(2-((2-95 mercaptophenyl)thio)ethyl)amine, L=CO (4), PMe₃ (5), NH₃ (6) and N₂H₄ (7)). Compounds 96 **4–5** have a LS (singlet) state and compounds **6–7** reportedly a HS (quintet) ground state. ^{46, 47,} 97 48 Fe^{III}(porphyrinato)Cl, Furthermore, focus on FePCl 98 we will (8) and, Fe^{III}(porphyrazinato)Cl, FePzCl (9), which have a sextet and a quartet ground state, 99 respectively. 100
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The optimization of the three Fe(III) molecules (1-3) leads in all cases to the expected structural changes for the different spin states (Tables S1-S3). Comparison of the optimized structures of 1-3 indicates the existence of an expansion of the ligand sphere. Going from the doublet to the quartet state, first the equatorial ligands move away from iron while the axial ligands stay almost at the same position. In the sextet state, the equatorial ligands remain virtually at the same position, but the axial ligands (have to) move out.

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Comparing the vertical spin state energies, calculated on the experimental structure,¹⁴ with 111 results from the optimized ("relaxed") geometries, Table 1, it is evident that the energy gap 112 between different spin states decreases. In the case of compound 1, the doublet state remains 113 the spin ground state with the quartet state (from 22.5 kcal·mol⁻¹ "vertical" to 17.5 kcal·mol⁻¹ 114 "relaxed") and the sextet state (from 33.9 kcal·mol⁻¹ "vertical" to 10.2 kcal·mol⁻¹ "relaxed") 115 in closer energetic proximity after geometry optimization. Molecule 2 has the quartet ground 116 state, and also here are the relative energies of the doublet and sextet states reduced after 117 structure relaxation. The same trends apply for the sextet ground state of complex **3**. For all 118 119 complexes, after the spin state relaxation, both OPBE and its recently developed successor, S12g, give the correct spin ground state. Spin contamination is small for these complexes, 120 121 and therefore shall not be discussed any further.

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123 The choice of exchange-correlation functional has an obvious influence on the geometry, with a tendency of S12g to give somewhat longer bond lengths than OPBE (Tables S1-S3). It 124 125 should be noted that S12g gives structural parameters that are in excellent agreement with experimental values. Unlike the choice of functional, the influence of solvation on the 126 geometrical parameters during the structural relaxation was not very significant, and it 127 depends slightly on system under consideration. In most cases, optimizations with COSMO 128 gave slightly longer bonds, but without significant consequences for the spin-state splittings, 129 Table I. 130

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<Table I>

- 134 Structure relaxation and spin state energies of compounds 4-7
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The spin state dependent structure relaxation for the Fe(II) compounds results in similar differences in Fe-ligand distances as for the Fe(III) compounds (Tables S4-S7). In the case of compounds **4-7**, the Fe-N, Fe-S and Fe-C distances are slightly elongated in comparison to the distances in Fe(III) complexes due to the additional *d*-electron in Fe(II) systems.

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The spin ground state of the Fe(II) complexes 4 and 5 is correctly predicted using both 141 OPBE and S12g levels of theory (see Table II): the singlet state is the lowest in energy for 142 both molecules, in agreement with experimental data. For compound 4, the triplet and quintet 143 144 states are significantly higher in energy. The energy differences between the different states are smaller for compound 5. Similar to the Fe(III) complexes, spin contamination is small and 145 shall not be discussed any further. Similar to compounds 4 and 5, we found after spin state 146 structure relaxation a LS ground state for iron complexes 6 and 7, with the IS and HS higher 147 in energy. Unfortunately, the experimental determination of the spin states of compounds $\mathbf{6}$ 148 149 and 7 is inconclusive, since anomalous high μ_{eff} values of 10-13 μ_B have been measured that may indicate impurities, e.g. by metallic iron, or oligomer formation. For compound 7 in 150 solution, a HS state was observed,⁴⁸ but a compound similar to 7 showed a diamagnetic LS 151 Fe center. ⁴⁹ Moreover, indications of dimer formation of the ligand-free [(Fe(NH)S₄)]-152 complex were observed. ^{47, 48} Since the different forms of the (Fe(NH)S₄)L complex in these 153 studies were obtained, we have checked both forms for compounds 4-7, i.e. with the "trans" 154 and "meso" form (see Figure 2). For both forms of each of compounds 4-7 we found a LS 155 156 ground state, albeit with smaller spin-state splitting for compounds 6 and 7. These findings can be traced back to the strength of the iron-ligand bond, which seems to be much weaker 157 for compounds 6/7 than for compounds 4/5. The weakly bound NH₃ and N₂H₄ ligands are 158 easily exchanged with CH₃OH, solvent (THF) or CO. ⁴⁸ These experimental data corroborate 159 our computed ligand-binding energies, which indicate strong and favorable binding of CO 160 and P(Me)₃ to form the singlet ground state, but less favorable binding of the other ligands or 161 162 spin states (see Table S10). Interestingly, the monomeric Fe(NH)S₄ complex without a ligand is predicted to have a triplet spin ground state in the trans form, with the other spin states or 163 the meso form lying higher in energy by at least 7 kcal·mol⁻¹. The ligand-free complex may 164 dimerize to give the experimentally observed HS state through ferromagnetic coupling. The 165 latter process has not been studied due to the complexity involved with ferromagnetic versus 166 anti-ferromagnetic coupling of the many spin states that need to be considered. This is 167 confirmed by a recent study using high-level *ab initio* methods that indeed found a singlet 168

ground-state for these molecules.⁵⁰ In another recent study, "accurate" spin ground states for 169 molecules 6 and 7 were found with the double hybrid B2PLYP functional, where a HS 170 ground-state was obtained for molecule 6 with OPBE.⁵¹ Since the last result is in 171 disagreement with our study, we re-optimized molecules 6 and 7 using the OPBE functional 172 with the geometries from their paper⁵¹ as starting point. The reoptimized structures resulted in 173 spin state splittings that were in accordance with the previous study,⁵¹ however, the structures 174 were highly distorted representing only a local minimum on the potential energy surface (and 175 ca. 5-20 kcal·mol⁻¹ above the structures obtained here for Table III). 176 177 <Figure 2> 178 179 As in the case of previous investigated molecules 1-3, after optimization with S12g, 180 somewhat longer bond lengths have been obtained in comparison with OPBE geometries. 181 Nevertheless both of them are again in good agreement with the experimental data (Tables 182 S4-S7). In contrast to the Fe(II) complexes (1-3), the Fe(III) P450 model systems (4-7) are 183 prone to the influence of solvent (COSMO calculations) on spin state ordering, Table II and 184 Table III. Calculations with COSMO solvation model revealed a tendency to favor the HS 185 186 state for complexes 5-7, and the LS state for complex 4. 187 188 <Table II> <Table III> 189 190 Iron porphyrin chloride and the porphyrazine analogue 191 192 The structures of FePCl (8) and FePzCl (9) were optimized in $C_{4\nu}$ symmetry, separately 193 for each spin state. Similarly to previous results,^{52, 53,54} it was found that the porphyrin core 194 size increases when going from the LS to the HS state, while the Fe-Cl distance increases 195 from the LS to IS state, and then slightly decreases in the HS state (Tables S8-S9). 196 OPBE and S12g predict the correct sextet spin ground state for both, FePCl and FePzCl 197 (see Table IV). In the case of FePCl a sextet ground state was predicted with the quartet 198 higher in energy and vice versa for FePzCl the quartet state is lower in energy. In both cases 199 200 the LS state is considerably higher in energy. 201 COSMO calculations revealed a clear and unambiguous solvent effect on the electronic 202

203 structure, Table IV. Introduction of the solvent favors the LS state, and as such have small effects on the spin ground state of molecule 9 that has a quartet ground state and a sextet 204 quartet state that is close in energy. In contrast, for molecule 8, that is in a HS experimentally 205 and has low lying quartet state, the quartet state is stabilized to the extent that it becomes the 206 ground state within all COSMO calculations. Of course, it should be added that we have been 207 investigating the spin-state splittings through looking at the electronic energy, and hence 208 ignoring the enthalpy and entropy effects. Both of these favor the high-spin states. Finally, 209 S12g once again shows excellent agreement with spin state energetics obtained with OPBE 210 211 level of theory.

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CONCLUSION

<Table IV>

Within this paper the extension of our previous validation¹⁴ of DFAs for a correct description of spin states of Fe(II) and Fe(III) complexes is presented. In the present contribution we allow the structure relaxation of the LS, IS and HS states of the iron compounds separately at OPBE and S12g levels of theory, thereby performing a more stringent test on the reliability of functionals for providing spin ground states of iron complexes.

A detailed comparison with the already proven OPBE DFA for spin state energetics, and experimental findings, revealed that S12g performs remarkably well, and thus represents a very promising tool for studying spin states in complicated transition metal systems. Moreover, for all iron complexes under investigation S12g gives a good match with experimental geometries, and thus can be considered as good starting point to investigate transition metal compounds.

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243	ИЗВОД
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245 246	РЕЛАКСАЦИЈА СПИНСКИХ СТАНЈА КОД КОМПЛЕКСА ГВОЖЂА: СЛУЧАЈ ЗА ОРВЕ И S12g ФУНКЦИОНАЛЕ
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248	МАЈА ГРУДЕН, ¹ СТЕПАН СТЕПАНОВИЋ ² and MARCEL SWART ^{*3,4}
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258	Структуре девет комплекса гвожђа који показују разноврсност експериментално
259	одређених основних спинских стања оптимизоване су Теоријом Функционала Густине
260	(DFT), а затим анализиране коришћењем различитих функционала. Извршена је
261	обимна валидациона студија новог S12g функционала, са дискусијом о утицају
262	окружења, геометрије, као и његових перформанси у односу на ОРВЕ функционал који
263	се већ показао као добар. ОРВЕ и S12g тачно предвиђају основно спинско стање код
264	свих испитиваних комплекса гвожђа. Како се S12g показао изузетно добро, он се може
265	сматрати поузданим за проучавање енергетике спинских стања у компликованим
266	системима прелазних метала.

268 Table captions:

- Table I Spin state energies (kcal mol⁻¹) for Fe(III) molecules 1-3 using TZP basis set,
 with OPBE and S12g functionals, in vacuum and COSMO.
- Table II Spin state energies (kcal mol⁻¹) for labile ("trans") complexes 4 and 5 using
 TZP basis, with OPBE and S12g functionals, in vacuum and COSMO.
- Table III Spin state energies (kcal mol⁻¹) for labile ("meso") complexes 6 and 7 using
 TZP basis, with OPBE and S12g functionals, in vacuum and COSMO.
- Table IV Spin state energy differences (kcal mol⁻¹, TZP basis) for FePCl (8) and
 FePzCl (9), with OPBE and S12g functionals, in vacuum and COSMO.

		Fe	e-(PyPepS) ₂ 1		Fe(tsalen)Cl 2			$Fe(N(CH_2-o-C_6H_4S)_3)(1MIm^b)$ 3		
Geo.*	SP [†]	doublet	quartet	sextet	doublet	quartet	sextet	doublet	quartet	sextet
	OPBE	0	17.1	10.2	6.5	0	3.9	6.6	7.9	0
	OPBE cosmo	0	19.4	13.0	9.3	0	6.9	7.9	7.4	0
OPBE	S12g	0	15.8	8.7	7.6	0	3.8	6.8	7.2	0
	S12g cosmo	0	18.2	11.6	10.2	0	6.4	8.2	6.8	0
	OPBE	0	18.8	13.1	5.2	0	2.9	6.2	7.5	0
ODDE	OPBE cosmo	0	17.4	10.2	9.7	0	7.7	8.0	7.2	0
cosmo	S12g	0	18.4	13.3	6.0	0	3.0	6.5	6.8	0
	S12g cosmo	0	17.1	10.6	10.2	0	7.4	8.3	6.5	0
	OPBE	0	18.3	10.5	7.4	0	6.2	7.6	8.1	0
S12 σ	OPBE cosmo	0	22.7	14.7	10	0	9.2	8.6	7.1	0
5125	S12g	0	15.4	8.7	7.5	0	6.6	6.5	7.0	0
	S12g cosmo	0	19.9	13.1	9.9	0	9.3	7.7	6.1	0
	OPBE	0	17.5	10.6	7.0	0	4.7	7.5	8.4	0
S12g	OPBE cosmo	0	20.5	14.9	11.2	0	6.7	8.8	7.2	0
cosmo	S12g	0	15.7	9.2	6.6	0	4.7	6.1	6.8	0
	S12g cosmo	0	18.8	13.7	10.4	0	6.5	7.6	5.9	0

Table I. Spin state energies (kcal mol⁻¹) for Fe(III) molecules 1-3 using TZP basis set, with
OPBE and S12g functionals, in vacuum and COSMO.

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^{*} Geometry optimization with frozen core electrones

⁺ Subsequent single point calculations with full electron basis sets

Table II. Spin state energies (kcal mol⁻¹) for labile ("trans") complexes **4** and **5** using TZP

basis, with OPBE and S12g functionals, in vacuum and COSMO.

Å		trans-(Fe(NH)S ₄)CO 4		<i>trans</i> -(Fe(NH)S ₄)PMe ₃ 5			
Geo. [‡]	SP§	singlet	triplet	quintet	singlet	triplet	quintet
	OPBE	0	23.4	34.8	0	16.3	20.1
ODDE	OPBE cosmo	0	24.5	36.6	0	17.3	18.6
OPBE	S12g	0	19.1	28.0	0	14.5	17.7
	S12g cosmo	0	20.2	29.7	0	15.4	16.3
	OPBE	0	23.5	35.3	0	16.4	20.4
ODDE	OPBE cosmo	0	24.5	36.5	0	17.3	18.3
cosmo	S12g	0	19.4	29.1	0	15.1	19.1
	S12g cosmo	0	20.3	30.3	0	15.9	17.1
	OPBE	0	23.4	34.2	0	19.6	19.4
\$12 σ	OPBE cosmo	0	24.3	36.4	0	20.3	19.2
012g	S12g	0	18.7	29.3	0	15.6	16.8
	S12g cosmo	0	19.6	31.4	0	16.3	16.6
	OPBE	0	24.6	35.0	0	19.9	19.2
S12g	OPBE cosmo	0	24.8	36.5	0	20.6	18.2
cosmo	S12g	0	20.4	30.8	0	15.8	17.1
	S12g cosmo	0	20.5	32.2	0	16.4	16.2

[‡] Geometry optimization with frozen core electrones

[§] Subsequent single point calculations with full electron basis sets

Table III. Spin state energies (kcal mol⁻¹) for labile ("meso") complexes 6 and 7 using TZP

basis, with OPBE and S12g functionals, in vacuum and COSMO.

			Fe(NH)S	4)NH3 6	meso-(Fe(NH)S ₄)N ₂ H ₄ 7			
Geo.**	SP ^{††}	singlet	triplet	quintet	singlet	triplet	quintet	
	OPBE	0	10.3	6.6	0	11.3	6.6	
	OPBE cosmo	0	10.1	3.9	0	10.7	4.4	
OPBE	S12g	0	7.7	2.6	0	8.5	2.5	
	S12g cosmo	0	7.4	-0.1	0	7.9	0.3	
	OPBE	0	10.6	7.2	0	11.5	7.1	
	OPBE cosmo	0	9.9	3.5	0	10.1	3.8	
COPBE	S12g	0	7.9	3.5	0	9.7	3.6	
	S12g cosmo	0	7.0	-0.2	0	8.2	0.5	
	OPBE	0	10.1	7.5	0	11.1	7.6	
S12-	OPBE cosmo	0	10.7	5.7	0	11.0	6.8	
512g	S12g	0	8.4	5.3	0	9.3	5.3	
	S12g cosmo	0	8.7	3.4	0	9.2	4.6	
	OPBE	0	10.1	6.8	0	10.8	6.6	
S12g	OPBE cosmo	0	10.5	4.4	0	10.4	5.2	
cosmo	S12g	0	8.7	4.9	0	9.4	4.6	
	S12g cosmo	0	8.8	2.4	0	8.9	3.2	

^{**} Geometry optimization with frozen core electrones** Subsequent single point calculations with full electron basis sets

288	Table IV. Spin state energy	y differences	(kcal mol ⁻¹ ,	, TZP basis) f	for FePCl (8)	and FePzCl (9),
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289	with OPBE and S1	2g functionals,	in vacuum	and COSMO.
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	GD ⁸⁸	FePCl FePzCl					
Geo.**	SP ³⁸	doublet	quartet	sextet	doublet	quartet	sextet
	OPBE	18.4	3.9	0	12.5	0	3.7
ODDE	OPBE cosmo	16.3	-1.0	0	15.6	0	7.6
OPBE	S12g	15.7	1.5	0	12.8	0	4.9
	S12g cosmo	13.8	-2.9	0	15.8	0	8.6
	OPBE	18.0	4.8	0	11.6	0	2.9
ODDE	OPBE cosmo	16.9	-1.7	0	16.6	0	8.2
cosmo	S12g	15.0	2.3	0	11.8	0	4.2
	S12g cosmo	14.1	-3.6	0	16.5	0	9.2
	OPBE	18.6	4.0	0	12.7	0	3.7
S12a	OPBE cosmo	16.3	-0.7	0	15.7	0	7.5
512g	S12g	15.4	1.5	0	12.6	0	5.0
	S12g cosmo	13.4	-2.8	0	15.4	0	8.6
	OPBE	18.3	4.5	0	12.0	0	3.2
S12g	OPBE cosmo	17.1	-1.6	0	16.8	0	8.1
cosmo	S12g	15.0	2.0	0	11.8	0	4.6
	S12g cosmo	14.0	-3.5	0	16.3	0	9.2

 ^{**} Geometry optimization with frozen core electrones
 ^{\$§} Subsequent single point calculations with full electron basis sets

292	Figure	captions:
293	٠	Figure 1. Fe-(PyPepS) ₂ 1 (PyPepSH ₂ = N-2-mercaptophenyl-2'-pyridine-
294		carboxamide); Fe(tsalen)Cl 2 (tsalen =N,N'-ethylenebis-(thio-salicylideneiminato));
295		Fe(N(CH ₂ -o-C ₆ H ₄ S) ₃)(1-Me-imidazole) 3 ; (Fe(NH)S ₄)L 4 (L=CO), 5 (PMe ₃), 6
296		(NH_3) , 7 (N_2H_4) ($(NH)S_4 = bis(2-((2-mercaptophenyl)thio)ethyl)amine); Iron$
297		porphyrin chloride (8, FePCl) and iron porphyrazine chloride (9, FePzCl).
298	٠	Figure 2. Different forms of compounds 4-7 .



301

Figure 1. Fe-(PyPepS)₂ **1** (PyPepSH₂ = N-2-mercaptophenyl-2'-pyridine-carboxamide); Fe(tsalen)Cl **2** (tsalen =N,N'-ethylenebis-(thio-salicylideneiminato)); Fe(N(CH₂-o-C₆H₄S)₃)(1-Me-imidazole) **3**; (Fe(NH)S₄)L **4** (L=CO), **5** (PMe₃), **6** (NH₃), **7** (N₂H₄) ((NH)S₄ = bis(2-((2-mercaptophenyl)thio)ethyl)amine); Iron porphyrin chloride (**8**, FePCl) and iron porphyrazine chloride (**9**, FePzCl).



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